

## COHERENT STRUCTURES IN A SELF-SIMILAR ADVERSE PRESSURE GRADIENT TURBULENT BOUNDARY LAYER AT THE VERGE OF SEPARATION

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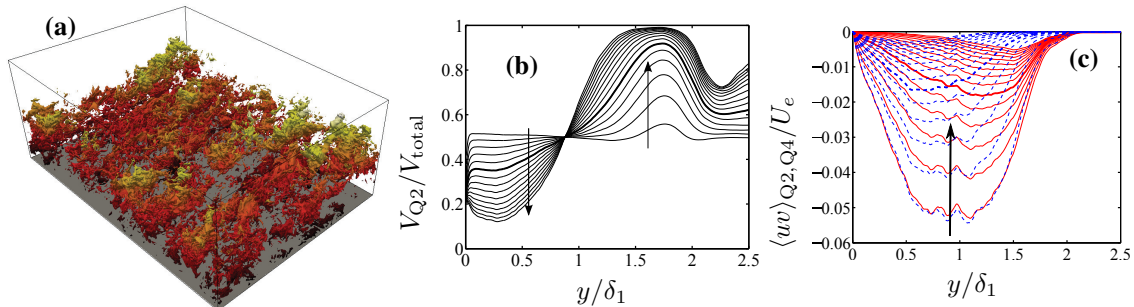
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The coherent structures of a self-similar adverse pressure gradient turbulent boundary layer (APG-TBL) at the verge of separation [7, 5] are investigated using a direct numerical simulation (DNS). The axes in the streamwise, wall-normal and spanwise directions are  $x$ ,  $y$  and  $z$ . The corresponding velocity fluctuation with respect to the time-averaged mean  $(U, V, W)$  are  $(u, v, w)$ . The APG-TBL achieves a region of constant friction coefficient, pressure velocity and shape factor [3] with turbulence statistics in this region showing a self-similar collapse by using the scaling of the external velocity  $(U_e)$  and the displacement thickness  $(\delta_1)$  and velocity fluctuations in this strong APG-TBL are mainly produced in the outer layer [4]. As shown in figure 1(a), the intense Reynolds stress structures ( $uv$ -structures) are detached from the wall and that is true for the fine-scale vortex clusters represented by isosurfaces of the second invariant of the velocity gradient tensor (not shown). Most of events are generated in the outer region of APG-TBL, and the interaction of the local shear and the Reynolds stress ( $uv$ -structures) are of great interest, in comparison with those in zero-pressure-gradient TBL, turbulent mixing layer and homogeneous shear turbulence (HST). Ejections (Q2) and sweeps (Q4) are extracted by  $\tau_{xy}^* < -Hu^*v^*$ , where  $\tau_{xy}^* = \langle uv \rangle / (U_e(x)^2)$ ,  $u^* = u/U_e(x)$ ,  $v^* = v/U_e(x)$  and  $H$  is a constant threshold. As shown in figure 1(b,c), the volume fraction and the scaled stress conditioned on the  $uv$ -structure indicate that ejections and sweeps are balanced at a symmetric point  $y/\delta_1 \approx 0.88$  independently of  $H$ . Ejections mainly interacts with the free stream at around  $y/\delta_1 \gtrsim 1.5$ , and on the other hand, sweeps have larger contributions on the Reynolds stress than ejections below the symmetric point. The wall effect is limited within  $y/\delta_1 \lesssim 0.5$  and the strong APG-TBL at the verge of separation behaves like a free shear flow. The Corrsin shear parameter [1],  $S^* \equiv (\partial U / \partial y) q^2 / \varepsilon$ , where  $q^2 \equiv u^2 + v^2 + w^2$  and  $\varepsilon$  is the dissipation rate, is  $S^* \approx 9$  in the outer layer of the strong APG-TBL in good agreements with both HST [6], turbulent mixing layer and at the top of the logarithmic layer of wall-bounded flows [2]. Further investigations on the coherent structures in the self-similar APG-TBL, in comparison with the other shear flows, will be presented.

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**Figure 1.** (a) The isosurfaces of intense Reynolds stress in the self-similar APG-TBL. Only the domain of interest (DoI)  $[L_{DoI}, 0.27L_y, L_z] = [6.4, 3.43, 8.8]\delta_1(x_{DoI})$  is shown.  $x_{DoI}$  denotes the beginning of DoI, where  $Re_{\delta_1}(x_{DoI}) \approx 25000$  for the strong APG-TBL, and  $L_{DoI} \times L_y \times L_z$  represents the dimensions of DoI. The isosurfaces are  $-uv/U_e^2 = 0.016$ , coloured by the distance from the wall. The flow is from left to top-right. (b) The volume fraction of ejections (Q2) and sweeps (Q4), conditioned by  $\tau_{xy}^* < -Hu^*v^*$ . The lines are  $H = [0.25 : 0.25 : 4.0]$ , and the thick line represents  $H = 1.75$ . The arrow indicates the increase of  $H$ . (c) the Reynolds stress conditioned on the  $uv$ -structures; (red solid) Q2; (blue dashed) Q4. The lines are the same with (b).

### References

- [1] S. Corrsin. Local isotropy in turbulent shear flow. *NACA Research Memo.*, **58B11**, 1958.
- [2] J. Jiménez. Near-wall turbulence. *Phys. Fluids*, **25**:101302, 2013.
- [3] V. Kitsios, C. Atkinson, J. A. Sillero, G. Borrell, A. G. Gungor, J. Jiménez, and J. Soria. Direct numerical simulation of a self-similar adverse pressure gradient turbulent boundary layer. *Int. J. Heat and Fluid flow*, **61**:129–136, 2016.
- [4] V. Kitsios, A. Sekimoto, C. Atkinson, J.A. Sillero, G. Borrell, A. Gul Gungor, J. Jiménez, and J. Soria. Direct numerical simulation of a self-similar adverse pressure gradient turbulent boundary layer at the verge of separation. (*submitted*).
- [5] G. L. Mellor and D. M. Gibson. Equilibrium turbulent boundary layers. *J. Fluid Mech.*, **24**:225–253, 1966.
- [6] A. Sekimoto, S. Dong, and J. Jiménez. Direct numerical simulation of statistically stationary and homogeneous shear turbulence and its relation to other shear flows. *Phys. Fluids*, **28**:035101, 2016.
- [7] A. A. Townsend. The development of turbulent boundary layers with negligible wall stress. *J. Fluid Mech.*, **8**(01):143–155, 1960.